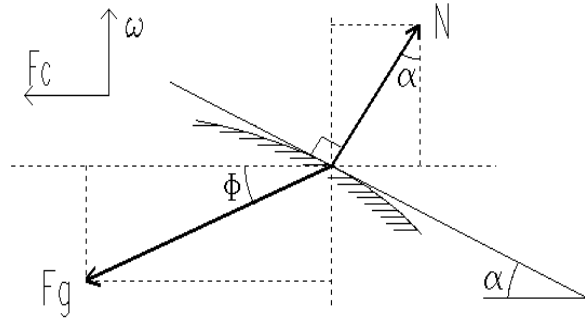


### Solution of question 3.

a - 1st method

For equilibrium we have  $F_c = F_g + N$   
where  $N$  is normal to the surface.

Resolving into horizontal and vertical components, we find:



$$F_g \cdot \cos(\phi) = F_c + N \cdot \sin(\alpha)$$

$$F_g \cdot \sin(\phi) = N \cdot \cos(\alpha) \rightarrow F_g \cdot \cos(\phi) = F_c + F_g \cdot \sin(\phi) \cdot \tan(\alpha)$$

From:

$$F_g = \frac{G \cdot M}{r^2}, \quad F_c = \omega^2 \cdot r, \quad x = r \cdot \cos(\phi), \quad y = r \cdot \sin(\phi) \text{ en } \tan(\alpha) = \frac{dy}{dx}$$

we find:

$$y \cdot dy + \left( 1 - \frac{\omega^2 \cdot r^3}{G \cdot M} \right) x \cdot dx = 0$$

where:

$$\frac{\omega^2 \cdot r^3}{G \cdot M} \approx 7 \cdot 10^{-4}$$

This means that, although  $r$  depends on  $x$  and  $y$ , the change in the factor in front of  $x dx$  is so slight that we can take it to be constant. The solution of Eq. (1) is then an ellipse:

$$\frac{x^2}{r_e^2} + \frac{y^2}{r_p^2} = 1 \rightarrow \frac{r_p}{r_e} = \sqrt{1 - \frac{\omega^2 \cdot r^3}{G \cdot M}} \approx 1 - \frac{\omega^2 \cdot r^3}{2 \cdot G \cdot M}$$

and from this it follows that:

$$\epsilon = \frac{r_e - r_p}{r_e} = \frac{\omega^2 \cdot r^3}{2 \cdot G \cdot M} \approx 3,7 \cdot 10^{-4}$$

2nd method

For a point mass of 1 kg on the surface,

$$U_{pot} = -\frac{G \cdot M}{r} \quad U_{kin} = \frac{1}{2} \cdot \omega^2 \cdot r^2 \cdot \cos^2(\phi)$$

The form of the surface is such that  $U_{pot} - U_{kin} = \text{constant}$ . For the equator ( $\Phi = 0$ ,  $r = r_e$ ) and for the pole ( $\Phi = \pi/2$ ,  $r = r_p$ ) we have:

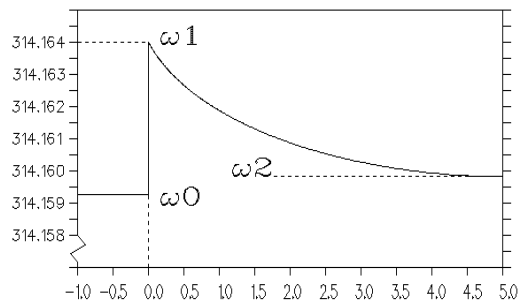
$$\frac{G \cdot M}{r_p} = \frac{G \cdot M}{r_e} + \frac{1}{2} \cdot \omega^2 \cdot r_e^2 \rightarrow \frac{r_e}{r_p} = 1 + \frac{\omega^2 \cdot r_e^3}{2 \cdot G \cdot M}$$

Thus:

$$\epsilon = \frac{r_e - r_p}{r_e} = \frac{1 + \frac{\omega^2 \cdot r_e^3}{2 \cdot G \cdot M} - 1}{1 + \frac{\omega^2 \cdot r_e^3}{2 \cdot G \cdot M}} \approx \frac{\omega^2 \cdot r_e^3}{2 \cdot G \cdot M} \approx 3,7 \cdot 10^{-4}$$

- b - As a consequence of the star-quake, the moment of inertia of the crust  $I_m$  decreases by  $\Delta I_m$ .

From the conservation of angular momentum, we have:



$$I_m \cdot \omega_0 = (I_m - \Delta I_m) \cdot \omega_1 \rightarrow \Delta I_m = I_m \cdot \frac{\omega_1 - \omega_0}{\omega_1}$$

After the internal friction has equalized the angular velocities of the crust and the core, we have:

$$(I_m + I_c) \cdot \omega_0 = (I_m + I_c - \Delta I_m) \cdot \omega_2 \rightarrow \Delta I_m = (I_m + I_c) \cdot \frac{\omega_2 - \omega_0}{\omega_2}$$

$$\frac{I_m}{I_m + I_c} = \frac{(\omega_2 - \omega_0) \cdot \omega_1}{(\omega_1 - \omega_0) \cdot \omega_2} \rightarrow 1 - \frac{I_c}{I_m + I_c} = \frac{(\omega_2 - \omega_0) \cdot \omega_1}{(\omega_1 - \omega_0) \cdot \omega_2}$$

$I (\cdot) R^2$

$$\rightarrow \frac{I_c}{I_m + I_c} = \frac{r_c^2}{r^2} \rightarrow \frac{r_c}{r} = \sqrt{1 - \frac{(\omega_2 - \omega_0) \cdot \omega_1}{(\omega_1 - \omega_0) \cdot \omega_2}} \approx 0.95$$

### Marking breakdown

- |   |            |   |      |
|---|------------|---|------|
| a | 1st method | - expressions for the forces                          | :1   |
|   |            | - equation for the surface                            | :2   |
|   |            | - equation of ellipse                                 | :1   |
|   |            | - flattening factor                                   | :1   |
|   | 2nd method | - energy equation                                     | :4   |
|   |            | - flattening factor                                   | :1   |
| b |            | - conservation of angular momentum for crust          | :1.5 |
|   |            | - conservation of angular momentum for crust and core | :1.5 |
|   |            | - moment of inertia for a sphere                      | :1   |
|   |            | - ratio $r_c/r$                                       | :1   |