

## Solutions to Theoretical Question 2

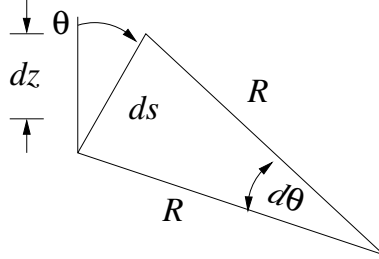
(a)

Snell's Law may be expressed as

$$\frac{\sin \theta}{\sin \theta_0} = \frac{c}{c_0} \quad , \quad (1)$$

where  $c$  is the speed of sound.

Consider some element of ray path  $ds$  and treat this as, locally, an arc of a circle of radius  $R$ . Note that  $R$  may take up any value between 0 and  $\infty$ . Consider a ray component which is initially directed upward from  $S$ .



In the diagram,  $ds = R d\theta$ , or  $\frac{ds}{d\theta} = R$ .

From equation (1), for a small change in speed  $dc$ ,

$$\cos \theta d\theta = \frac{\sin \theta_0}{c_0} dc$$

For the upwardly directed ray  $c = c_0 + bz$  so  $dc = b dz$  and

$$\frac{\sin \theta_0}{c_0} b dz = \cos \theta d\theta \quad , \quad \text{hence} \quad dz = \frac{c_0}{\sin \theta_0} \frac{1}{b} \cos \theta d\theta \quad .$$

We may also write (here treating  $ds$  as straight)  $dz = ds \cos \theta$ . So

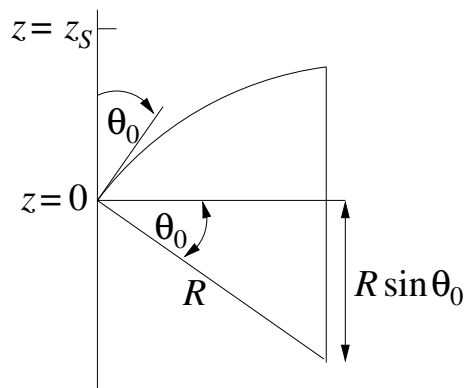
$$ds = \frac{c_0}{\sin \theta_0} \frac{1}{b} d\theta$$

Hence

$$\frac{ds}{d\theta} = R = \frac{c_0}{\sin \theta_0} \frac{1}{b} \quad .$$

This result strictly applies to the small arc segments  $ds$ . Note that from equation (1), however, it also applies for all  $\theta$ , i.e. for all points along the trajectory, which therefore forms an arc of a circle with radius  $R$  until the ray enters the region  $z < 0$ .

(b)



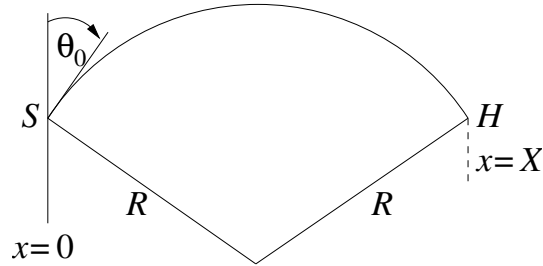
Here

$$\begin{aligned} z_s &= R - R \sin \theta_0 \\ &= R(1 - \sin \theta_0) \\ &= \frac{c_0}{b \sin \theta_0} (1 - \sin \theta_0) , \end{aligned}$$

from which

$$\theta_0 = \sin^{-1} \left[ \frac{c_0}{bz_s + c_0} \right] .$$

(c)



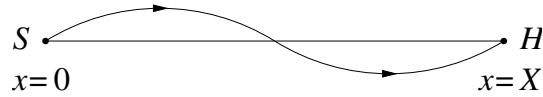
The simplest pathway between  $S$  and  $H$  is a single arc of a circle passing through  $S$  and  $H$ . For this pathway:

$$X = 2R \cos \theta_0 = \frac{2c_0 \cos \theta_0}{b \sin \theta_0} = \frac{2c_0}{b} \cot \theta_0 .$$

Hence

$$\cot \theta_0 = \frac{bX}{2c_0} .$$

The next possibility consists of two circular arcs linked as shown.



For this pathway:

$$\frac{X}{2} = 2R \cos \theta_0 = \frac{2c_0}{b} \cot \theta_0 .$$

i.e.

$$\cot \theta_0 = \frac{bX}{4c_0} .$$

In general, for values of  $\theta_0 < \frac{\pi}{2}$ , rays emerging from  $S$  will reach  $H$  in  $n$  arcs for launch angles given by

$$\theta_0 = \cot^{-1} \left[ \frac{bX}{2nc_0} \right] = \tan^{-1} \left[ \frac{2nc_0}{bX} \right]$$

where  $n = 1, 2, 3, 4, \dots$

Note that when  $n = \infty$ ,  $\theta_0 = \frac{\pi}{2}$  as expected for the axial ray.

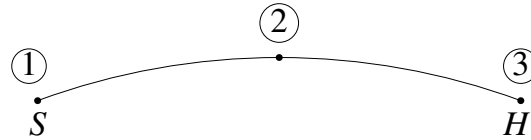
(d)

With the values cited, the four smallest values of launch angle are

$n$	$\theta_0$ (degrees)
1	86.19
2	88.09
3	88.73
4	89.04

(e)

The ray path associated with the smallest launch angle consists of a single arc as shown:



We seek

$$\int_1^3 dt = \int_1^3 \frac{ds}{c}$$

Try first:

$$t_{12} = \int_1^2 \frac{ds}{c} = \int_{\theta_0}^{\pi/2} \frac{Rd\theta}{c}$$

Using

$$R = \frac{c}{b \sin \theta}$$

gives

$$t_{12} = \frac{1}{b} \int_{\theta_0}^{\pi/2} \frac{d\theta}{\sin \theta}$$

so that

$$t_{12} = \frac{1}{b} \left[ \ln \tan \frac{\theta}{2} \right]_{\theta_0}^{\pi/2} = -\frac{1}{b} \ln \tan \frac{\theta_0}{2}$$

Noting that  $t_{13} = 2t_{12}$  gives

$$t_{13} = -\frac{2}{b} \ln \tan \frac{\theta_0}{2} .$$

For the specified  $b$ , this gives a transit time for the smallest value of launch angle cited in the answer to part (d), of

$$t_{13} = 6.6546 \text{ s}$$

The axial ray will have travel time given by

$$t = \frac{X}{c_0}$$

For the conditions given,

$$t_{13} = 6.6666 \text{ s}$$

thus this axial ray travels slower than the example cited for  $n = 1$ , thus the  $n = 1$  ray will arrive first.