

## Part 2a

The total energy radiated per second =  $4\pi R^2 \sigma T^4$ , where  $\sigma$  is the Stephan-Boltzmann constant. The energy incident on a unit area on earth per second is;

$$P = \frac{4\pi R^2 \sigma T^4}{4\pi \ell^2} \text{ yielding, } R = \left( P / \sigma T^4 \right)^{1/2} \ell \quad (1) \quad (0.8 \text{ pts})$$

The energy of a photon is  $hf = hc/\lambda$ . The equivalent mass of a photon is  $h/c\lambda$ . Conservation of photon energy:

$$\frac{hc}{\lambda_0} - \frac{Gm_0}{R} \cdot \frac{h}{c\lambda_0} = \frac{hc}{\lambda} \quad (0.8 \text{ pts})$$

yielding

$$R = \frac{Gm_0(\lambda_0 + \Delta\lambda)}{c^2 \Delta\lambda} \quad (2)$$

and (2) yields,

$$m_0 = \frac{c^2 \Delta\lambda \left( P / \sigma T^4 \right)^{1/2}}{G(\lambda_0 + \Delta\lambda)} \ell \quad (3) \quad (0.2 \text{ pts})$$

The stars are rotating around the center of mass with equal angular speeds:

$$\omega = (2\pi/2\tau) = \pi/\tau \quad (4) \quad (0.2 \text{ pts})$$

The equilibrium conditions for the stars are;

$$\frac{GMm_0}{(r_1 + r_2)^2} = m_0 r_1 \omega^2 = M r_2 \omega^2 \quad (5) \quad (0.8 \text{ pts})$$

with

$$r_1 = \ell \frac{\Delta\theta}{2}, \quad r_2 = \ell \frac{\Delta\phi}{2} \quad (6) \quad (0.4 \text{ pts})$$

Substituting (3), (4) and (6) into (5) yields

$$\ell = \left( \frac{8c^2 \Delta\lambda \left( P / \sigma T^4 \right)^{1/2}}{\Delta\phi(\pi/\tau)^2 (\lambda_0 + \Delta\lambda) (\Delta\theta + \Delta\phi)^2} \right)^{1/2} \cdot \quad (0.8 \text{ pts})$$

## Part 2b

Conservation of angular momentum for the ordinary star;

$$m r^2 \omega = m_0 r_0^2 \omega_0 \quad (7) \quad (0.6 \text{ pts.})$$

Conservation of angular momentum for dm:

$$r^2 \omega dm = r_f^2 \omega_f dm \quad (8) \quad (0.6 \text{ pts})$$

where  $\omega_f$  is the angular velocity of the ring. Equilibrium in the original state yields,

$$\omega_0 = \left( \frac{GM}{r_0^3} \right)^{1/2} \quad (9) \quad (0.8 \text{ pts})$$

and (7), (8) and (9) give,

$$\omega = \frac{m_0 r_0}{m r^2} \left( \frac{GM}{r_0} \right)^{1/2}, \quad \omega_f = \frac{m_0 r_0}{m r_f^2} \left( \frac{GM}{r_0} \right)^{1/2} \quad (10) \quad (0.4 \text{ pts})$$

Conservation of energy for dm;

$$\frac{1}{2} dm \left( v_0^2 + r^2 \omega^2 \right) - \frac{GM dm}{r} = \frac{1}{2} dm r_f^2 \omega_f^2 - \frac{GM dm}{r_f} \quad (11) \quad (1.2 \text{ pts})$$

Substituting (10);

$$v_0^2 + \frac{m_0^2 r_0 GM}{m^2} \left( \frac{1}{r^2} - \frac{1}{r_f^2} \right) - 2GM \left( \frac{1}{r} - \frac{1}{r_f} \right) = 0 \quad (12)$$

Since  $r_0 \gg r_f$ , if  $r > r_0$ ,  $r^{-1}$  and  $r^{-2}$  terms can be neglected. Hence,

$$r_f = \frac{GM}{v_0^2} \left( \left( 1 + \frac{m_0^2 r_0 v_0^2}{GM m^2} \right)^{1/2} - 1 \right). \quad (0.8 \text{ pts})$$

To show that  $r > r_0$  change in the linear momentum of the ordinary star in its reference frame:

$$-\frac{GMm}{r^2} + m r \omega^2 - m \frac{dv_r}{dt} = -v_0 \frac{dm_{gas}}{dt} \quad (13) \quad (0.8 \text{ pts})$$

and (13) implies the existence of an outward force initially and hence  $r$  starts growing. Using (7) one can write

$$m r \omega^2 = \frac{m_0^2 r_0^4 \omega_0^2}{m r^3}.$$

Hence,  $\frac{\text{Gravitational force}}{\text{Centrifugal force}} \propto m^2 r$ . (0.4 pts)

where  $m$  is definitely decreasing. If  $r$  starts decreasing at some time also, this ratio starts decreasing, which is a contradiction.

So  $r > r_0$ . (0.4 pts)